



PEGASUS-III
Experiment

Identifying Beam-Driven Instabilities Responsible for Current Drive on PEGASUS-III

R.K. Sassella

S.J. Diem, M.D. Nornberg, J.A. Reusch, A.T. Rhodes, C.E. Schaefer

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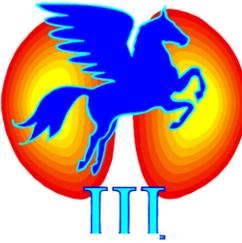
Presentation CP11.00048

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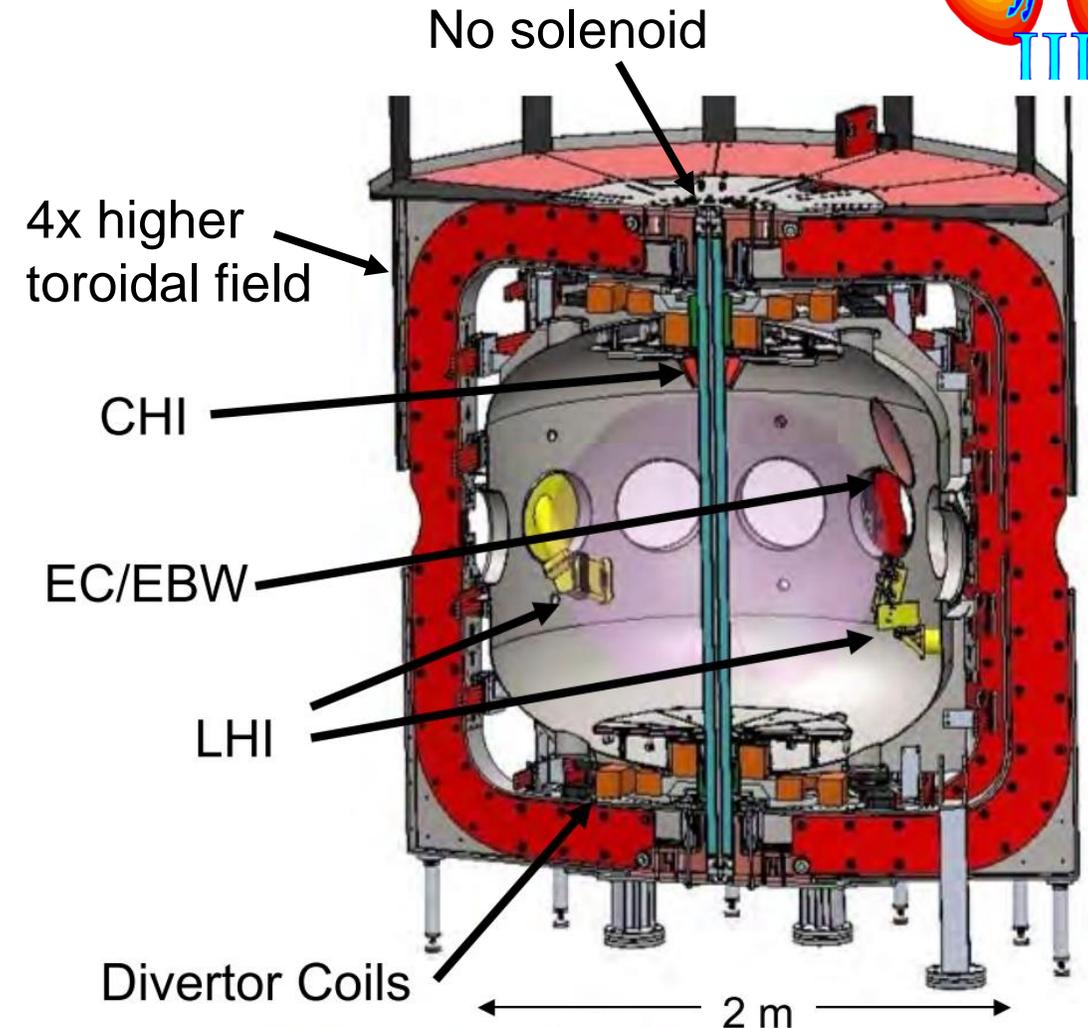
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PEGASUS-III Will Explore Non-Solenoidal Current Drive Techniques At Higher Toroidal Field



- Goal: testing and projecting performance of non-solenoidal startup to fusion plant scale devices
- Understanding **physical mechanisms behind current drive** crucial to extrapolate performance

Parameter	PEGASUS	PEGASUS-III
$I_{p, target}$	225 kA	300 kA
$B_{T, max}$	0.15 T	0.58 T
B_T Flattop	25 ms	50-100 ms
A	1.15	1.18

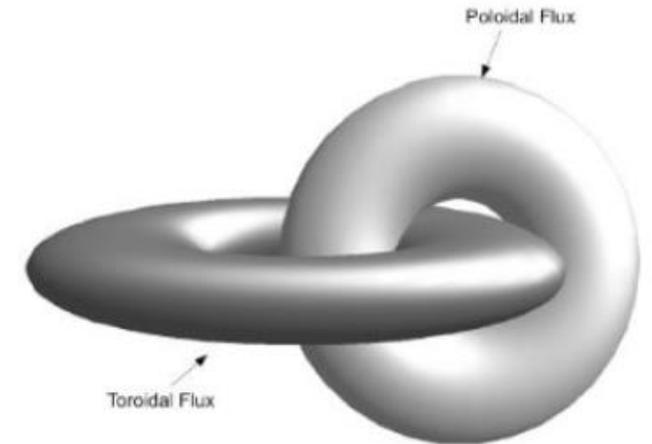


Helicity Injection Exploits Conservation Properties to Drive Current



- **Magnetic helicity** describes “interlinkedness” of flux in a volume
- In a tokamak, K_{plasma} results from linkage between Ψ_{poloidal} and Ψ_{toroidal}

$$K = \int \vec{A} \cdot \vec{B} d\tau$$
$$\propto I_p \Psi_{\text{tor}}$$

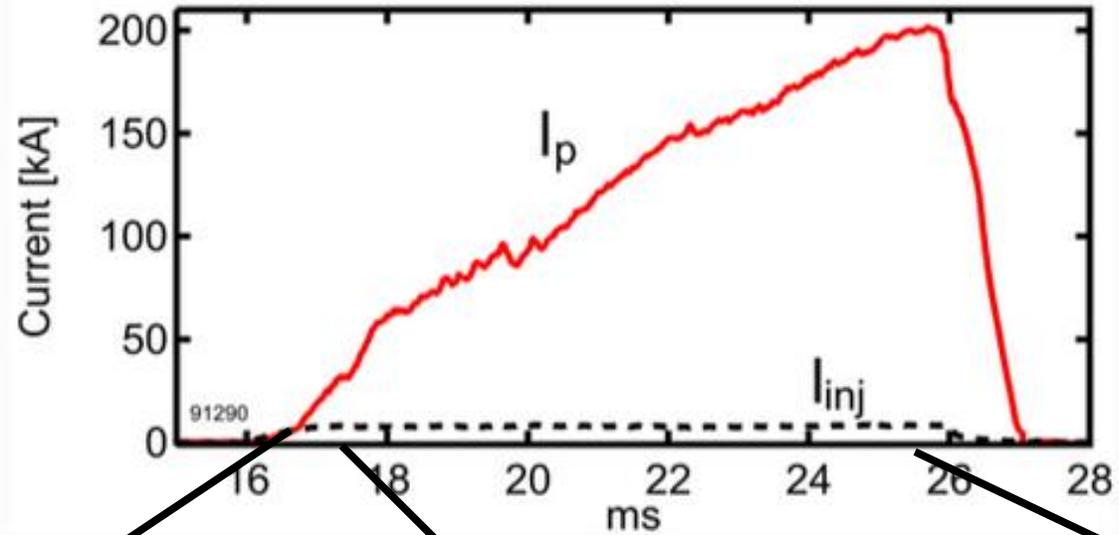


- Local helicity injection (LHI): biased cathode produces beam of electrons that spirals around vessel and relaxes into tokamak-like configuration
- Plasma current sustained by LHI proportional to injector properties, field at injector location:

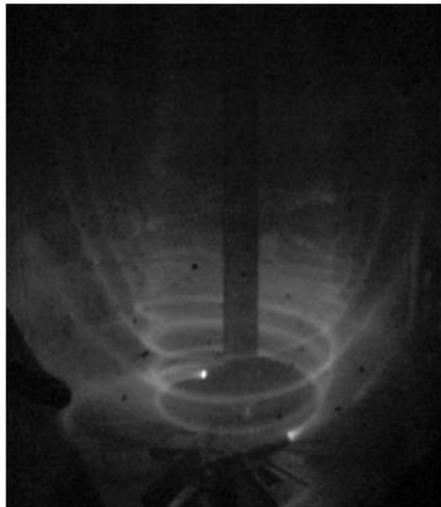
$$I_p = \frac{A_p V_{LHI}}{2\pi R_0 \langle \eta \rangle} \approx \frac{V_{inj} A_{inj}}{2\pi R_0 \langle \eta \rangle} \frac{B_{inj}}{B_0}$$



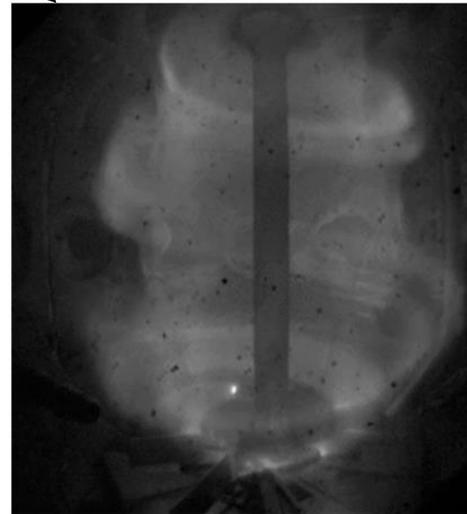
Relaxation Takes Plasma From LHI Startup To High Current Tokamak State



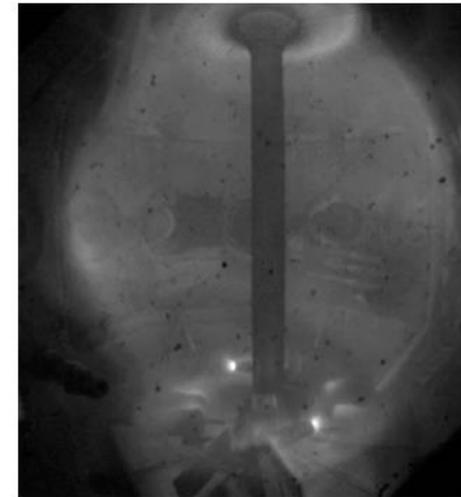
Slide adapted from: Richner, N. J.
"Observations of Magnetic
Turbulence During Local Helicity
Injection on PEGASUS." Invited
talk, APS DPP, November 2021.



$$I_p \approx N_{turns} I_{inj}$$



$$I_p \gtrsim N_{turns} I_{inj}$$



$$I_p \gg N_{turns} I_{inj}$$



Exploring Physical Mechanisms for LHI Current Drive

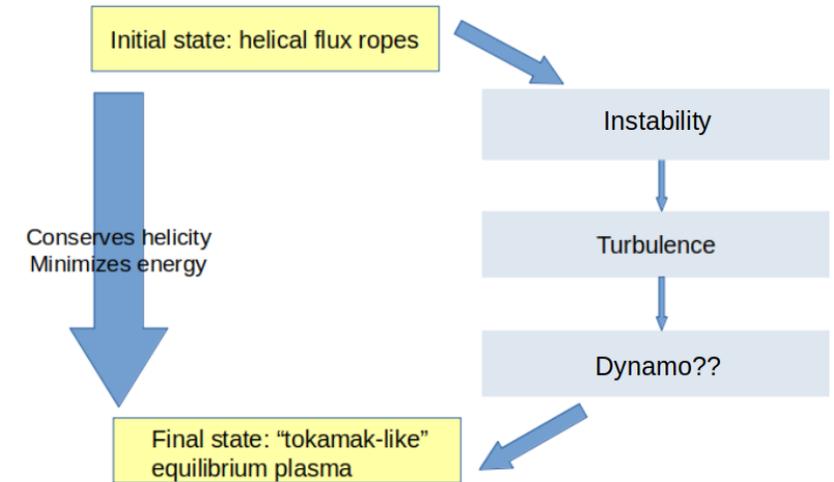
- Global limits for I_p are well understood
 - Helicity balance

$$I_p = \frac{A_p V_{LHI}}{2\pi R_0 \langle \eta \rangle} \approx \frac{V_{inj} A_{inj}}{2\pi R_0 \langle \eta \rangle} \frac{B_{inj}}{B_0}$$

- Taylor limit: I_p limited by edge $J_{||}/B$

$$I_p \leq \frac{\Psi_{tor} I_{inj}}{2\pi R_{inj} B_{p,inj} w}$$

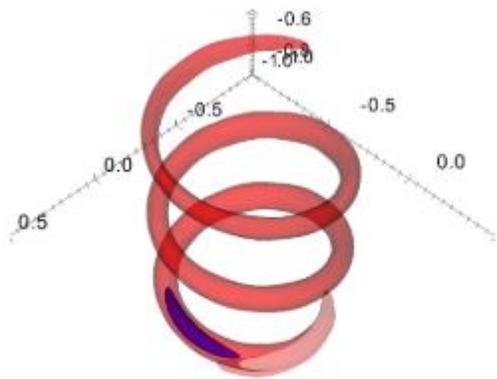
- Understanding **physical processes behind current drive** crucial to extrapolate LHI performance to fusion-relevant parameter spaces



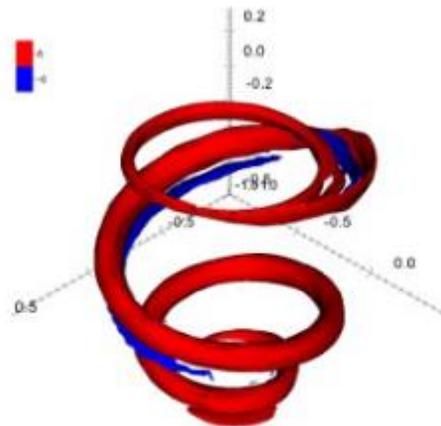
Simulations Predict a Relaxation Mechanism Based On Global Instability



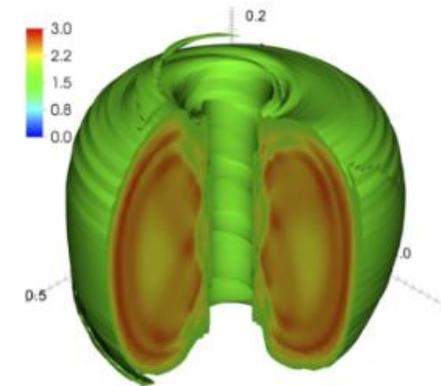
- NIMROD simulations model relaxation during LHI on PEGASUS (O'Bryan, J. B., & Sovinec, C. R. (2014). "Simulated flux-rope evolution during non-inductive startup in Pegasus." *Plasma Physics and Controlled Fusion*, 56(6).)
- Injected electron beams modeled as **flux ropes**
- Adjacent passes of flux rope reconnect and release current rings, which diffuse into tokamak-like plasma
- Simulations suggest a mechanism for poloidal flux development: the **dynamo**



initial current path

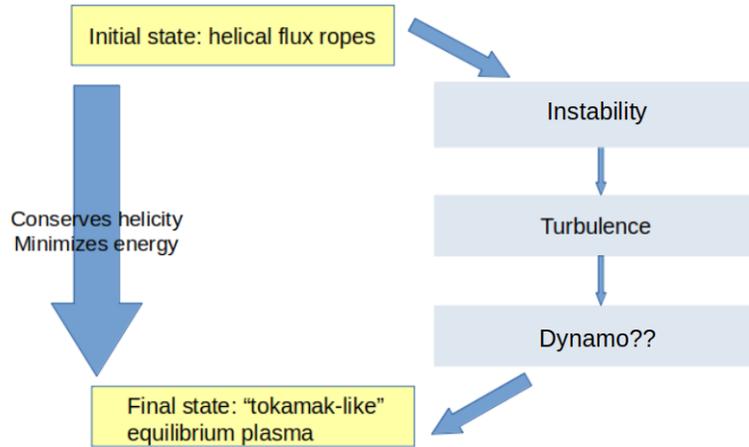


early ring formation



tokamak-like
axisymmetric plasma

Dynamo Effects Generate Large-Scale Fields and Flows



- Astrophysical dynamos convert kinetic energy to magnetic energy (flow-dominated)
 - Earth's magnetosphere
 - Galaxy-scale magnetic fields
- Lab plasma dynamos tend to be magnetically dominated
 - Dynamo current drive effects studied in depth on MST

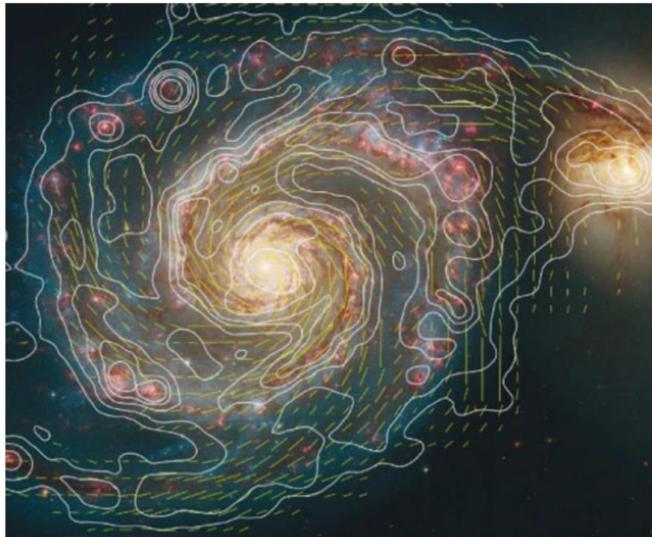


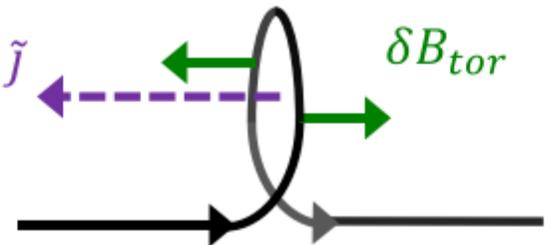
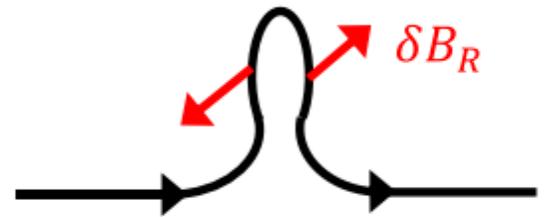
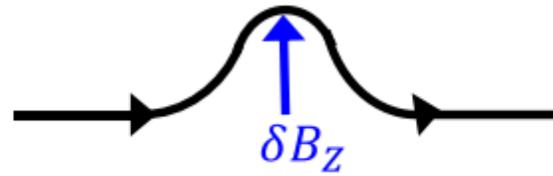
Image source: Rincon, F. (2019). Dynamo theories. Journal of Plasma Physics, 85(4). <https://doi.org/10.1017/s0022377819000539>



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Dynamos Can Drive Field-Parallel Current



- Lab dynamos typically involve fluctuation in magnetic field and another field quantity
- Example: in MHD dynamo, perpendicular velocity and magnetic fluctuations drive parallel current

$$\bar{\mathbf{E}} + \langle \tilde{\mathbf{v}} \times \tilde{\mathbf{B}} \rangle = \eta \bar{\mathbf{j}}$$

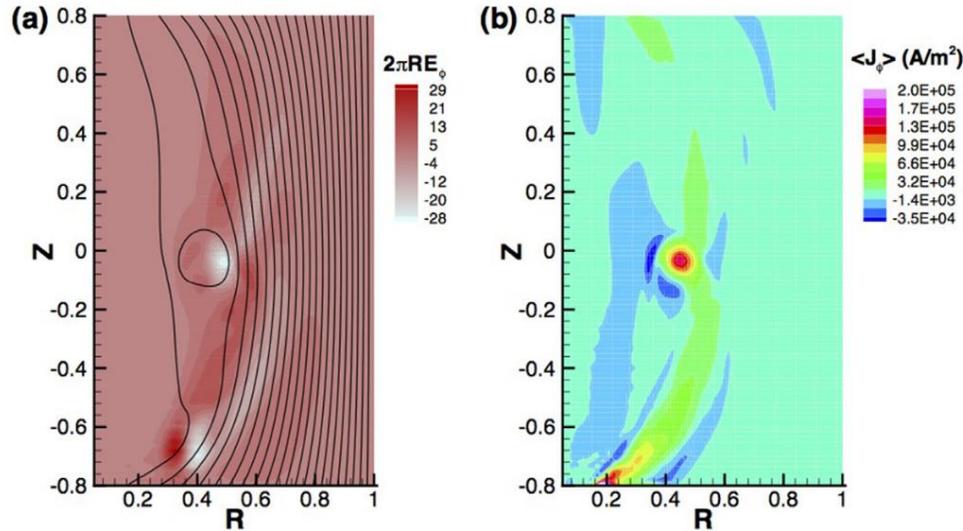
- Phase difference of fluctuations affects amount of current driven and its direction
- Dynamo activity must be considered statistically

Image source: Richner, N. J. "Observations of Magnetic Turbulence During Local Helicity Injection on PEGASUS." Invited talk, APS DPP, November 2021.



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Simulations Predict a Dynamo Mechanism Based On Global-Scale Instability on PEGASUS



Dynamo activity during an early ring formation event

a) Contours of induced EMF from MHD fluctuations

b) Toroidally averaged poloidal current density

Source: O'Bryan, J. B., & Sovinec, C. R. (2014). "Simulated flux-rope evolution during non-inductive startup in Pegasus." *Plasma Physics and Controlled Fusion*, 56(6).

Simulation predicted dynamo drive from flux-rope scale fluctuations from an **island coalescence instability**

\tilde{v} : slight rotation of flux rope

$\tilde{\mathbf{B}}$: inherently asymmetric field of flux rope



$$\bar{\mathbf{E}} + \langle \tilde{\mathbf{v}} \times \tilde{\mathbf{B}} \rangle = \eta \bar{\mathbf{j}}$$

→ Predicts dynamo current drive associated with n=1 mode on the order of 1E4 Hz



Observation of $n=1$ Mode During LHI Shows Global Modes Can't Fully Explain Current Drive

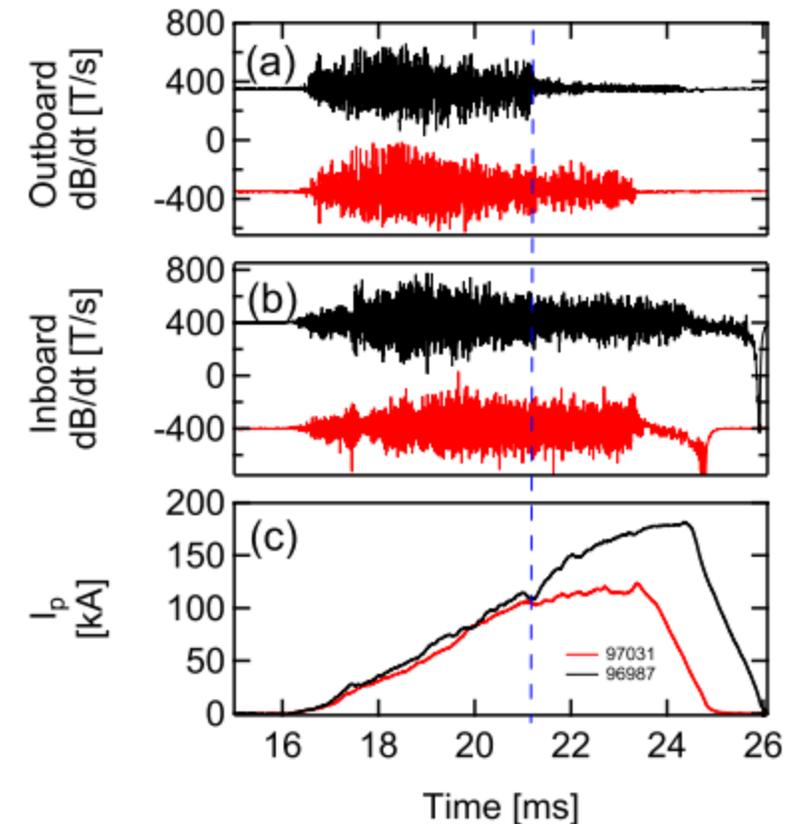


Observation: High-amplitude $n=1$ kink mode at 20 – 50 kHz dominates in a large parameter space

- Kink mode could result from island coalescence instability
- Suggested $n=1$ mode would be primary current drive mechanism, until...

A mode with $n=1$ activity reduced by $\sim 10\times$ (“**reduced MHD mode**”) was discovered on PEGASUS for HFS injection

- Kink reduction event corresponds to $\sim 50\%$ current drive *increase*
- High frequency fluctuations inside edge suddenly increase \rightarrow related to current drive?
- Motivated experiments to characterize magnetic turbulence and its relationship to current drive on PEGASUS



Source: Perry, J. M. et al. (2018). “Initiation and sustainment of tokamak plasmas with local helicity injection as the majority current drive.” *Nuclear Fusion*, 58(9).



LHI Generates Broadband Turbulence That Obeys Power Laws

- Turbulence characterized by uniform transfer (“**cascade**”) of energy to smaller scales at constant rate down to dissipation
- Helicity injected along with energy, but cascades up to system scale—develops large scale structures
- Turbulence in magnetized plasmas often **Alfvénic**: travels along field lines w/ \tilde{B} and \tilde{v} perpendicular to direction of propagation

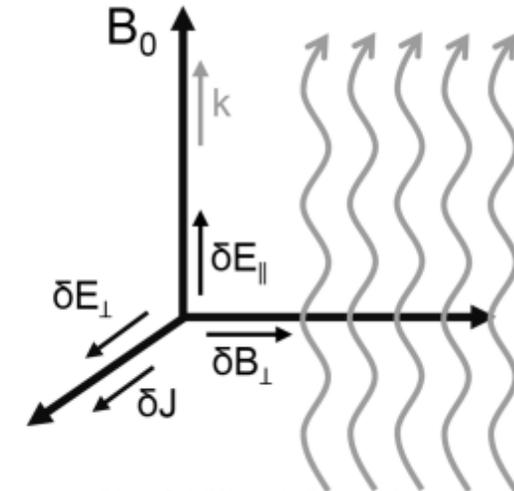
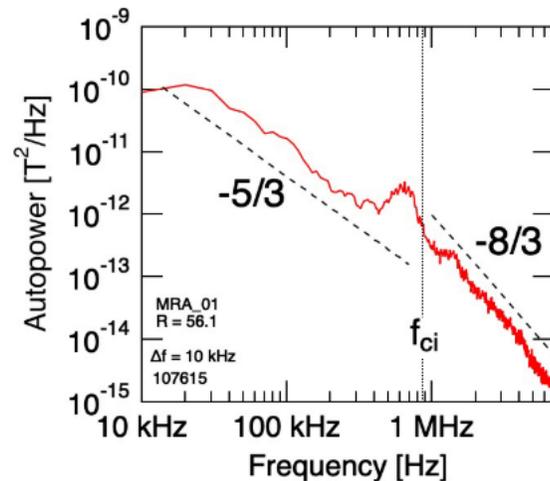


Image source: Diem, S.J. “Waves in Plasmas.” Introduction to Fusion and Plasma Physics Course lecture slides, SULI, 16 June 2020.



- Magnetic fluctuations measured on PEGASUS with a radial probe
- Measured LHI spectrum displays power law activity
 - MHD Alfvén waves at low frequencies (index = -5/3)
 - kinetic Alfvén waves at high frequencies (index = -8/3)

Source: Richner, N. J. “Observations of Magnetic Turbulence During Local Helicity Injection on PEGASUS.” Invited talk, APS DPP, November 2021.



Current is Driven Throughout the Alfvénic Cascade

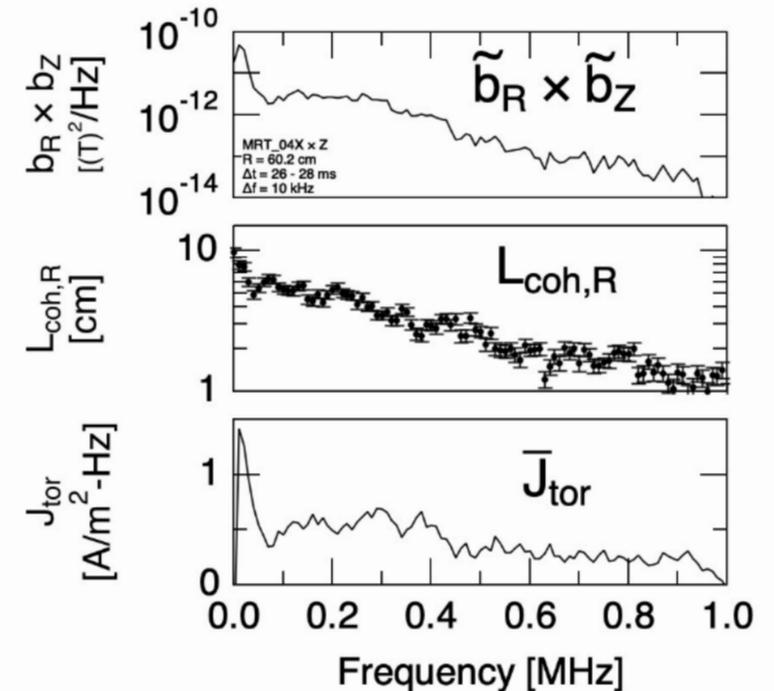
- Dynamo current can be estimated with the alpha dynamo model and Ampere's Law
- Correlations of fluctuations in b_z and b_r suggest current drive direction matches that of PEGASUS
- High frequency ($f \gg f_{ci}$) fluctuations contain significant power and correlate with total I_p
- Dynamo-driven current density estimate ~ 400 kA/m²; equilibrium reconstruction current $\sim 250 - 600$ kA/m²

→ Broadband (including $f \gg f_{ci}$) dynamo current drive may be a significant contributor to I_p on PEGASUS

$$\bar{J}_{tor}(f) \approx 2 \text{sgn}(\theta_{RZ}) \frac{\sqrt{\langle \tilde{b}_z \tilde{b}_R \rangle}}{\mu_0 L_{coh,R}}$$

Net coherent power

Co- vs. counter- drive



Future Work & Planned Experiments



Characteristics of turbulence parallel to magnetic field

- Multiple probes spaced out along a field line
- **Parallel correlation length:** distance over which turbulence decorrelates from itself; may be related to scale of instabilities in LHI

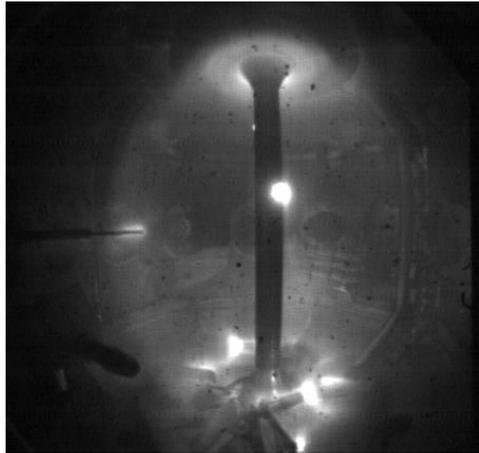
Scaling of current drive mechanism up to power plant scale

- Measure amplitude of turbulence at high field (up to 0.6 T)
- Extrapolate turbulent behavior, instability performance
- Requires better understanding of instability

Existing Magnetic Turbulence Probes Can Be Adapted for New Experiments



Probe Inserted into LHI Discharge

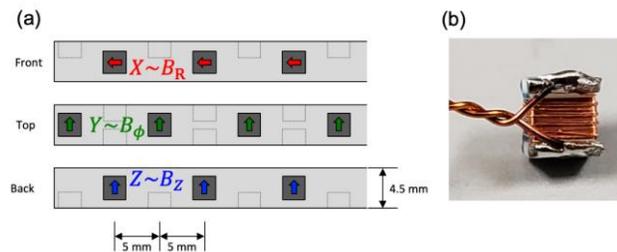


Magnetic probes built to measure magnetic fluctuations on PEGASUS (N. J. Richner)

- Features: frequency range up to MHz; variable insertion depth with a 40 cm throw; electrostatic and heat shielding
- Broadband array
- **Triple axis array**
 - 10 radially spaced, identical coil “triplets”
 - Co-located measurements of B_r , B_z , and B_ϕ

Specifications for new probe system:

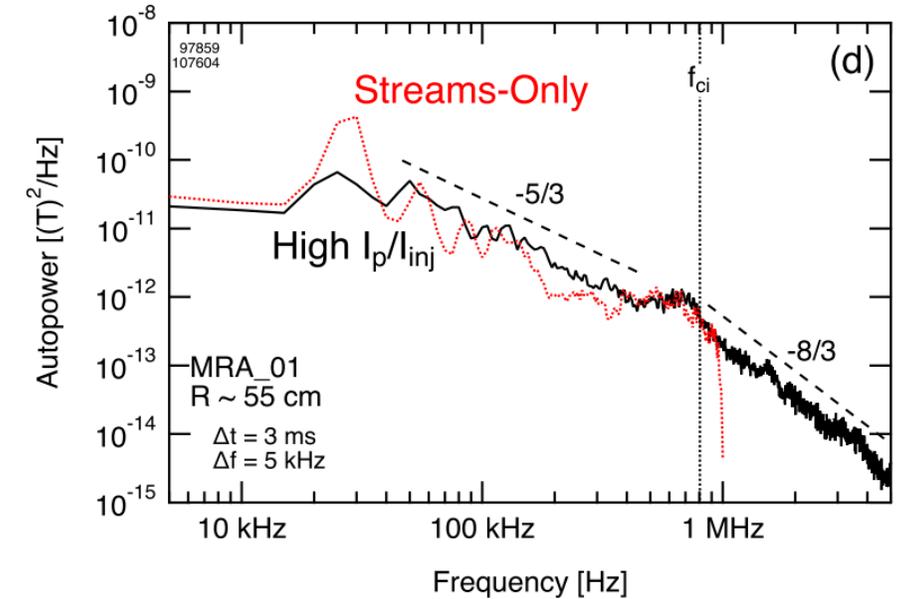
	bare minimum	desirable	aspirational
Dimensions	1 (radial)	3	3
Number of probes	2	3	4
Frequency range	2 MHz	5 MHz	10 MHz
Spatial resolution	<2 cm	<1 cm	< 5 mm
Max heating 10 ms	2500 K	1000 K	500 K





Future Work: Location and Type of Instability

- Evidence suggests that high- ω turbulence is generated by an injector or beam instability
 - Fluctuation amplitude increases with beam velocity
 - Streams-only discharges had similar spectra to relaxed discharges \rightarrow turbulence comes from source other than reconnection/relaxation process
- Possible instability sources:
 - Super-Alfvénic, super-thermal beam electrons streaming into background plasma induce waves all along beam length
 - “Double-layer” plasma structure around injector electrode known to have own set of instabilities



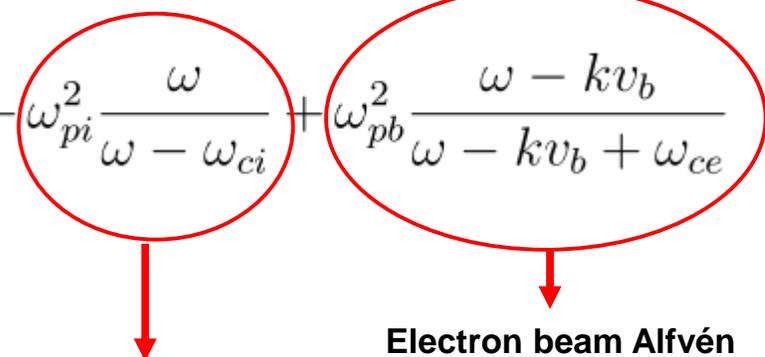
Source: Richner, N. J. (2021). Characterization of Magnetic Activity During Local Helicity Injection. University of Wisconsin-Madison.



Future Work: Instability Identification

- Modeling PEGASUS as monoenergetic beam incident on cold magnetized plasma, two possible instabilities:

$$D(\omega, k) = k^2 c^2 - \omega^2 + \omega_{pe}^2 \frac{\omega}{\omega + \omega_{ce}} + \omega_{pi}^2 \frac{\omega}{\omega - \omega_{ci}} + \omega_{pb}^2 \frac{\omega - kv_b}{\omega - kv_b + \omega_{ce}}$$



Electron beam ion cyclotron instability

Electron beam Alfvén wave instability (low ω , wave number limit)

- Modeling in kinetic regime (high k , ω limit) with different electron and ion temps predicts **kinetic Alfvén wave instability**

$$\omega^2 = k_{\parallel}^2 v_A^2 \left(\frac{k_{\perp}^2 \rho_i^2}{1 - I_0(k_{\perp}^2 \rho_i^2) \exp(-k_{\perp}^2 \rho_i^2)} + \frac{T_e}{T_i} k_{\perp}^2 \rho_i^2 \right)$$

- Calculating most unstable mode and growth rate as functions of parameter sweeps will enable instability ID